

# Critical chaos

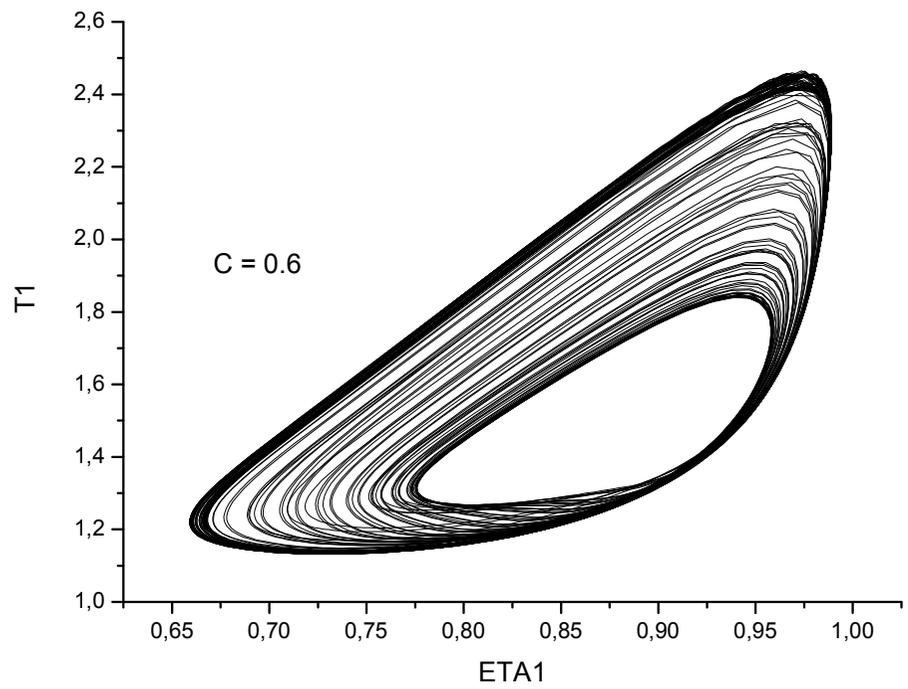
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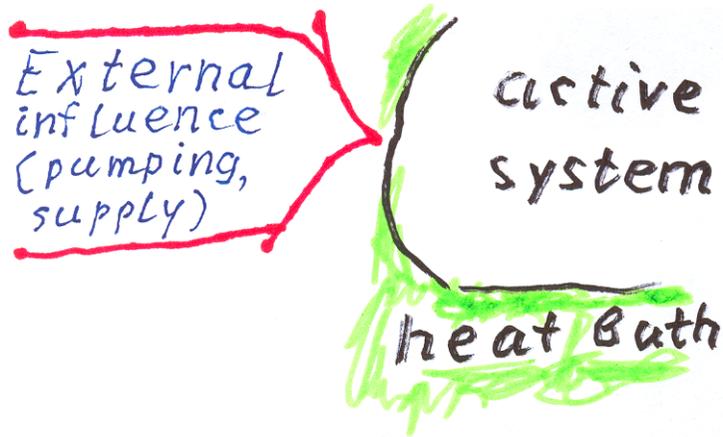
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There are two species of chaos, critical and dynamical. Trajectories of dynamical chaos are locally smooth, while for the critical kind they are of Brownian type. Critical chaos exists near bifurcations where unavoidable noise is amplified. The pulsating spectrum is independent of initial noise, soft modes being prevail. Onset of soft modes betokens catastrophe.

The following example is illustrating. Analytical results show that in the Edward Lorenz set with parameters close to the values of  $\theta = 8/3$ ,  $\sigma = 10.2$ ,  $r = 30.2$  (chaos boundary) stable limit cycles should exist. However, the cycles were not obtained by numerical simulation. Moreover, there always exists such a parameter region where, according to a rigorous theory, stable cycles can occur, and yet they escape detection by numerical means, whatever the accuracy of numerical analysis.

Such a failure is due to the noise. Any weak noise becomes strong as parameters approach the boundary, i.e. bifurcation.





space of states



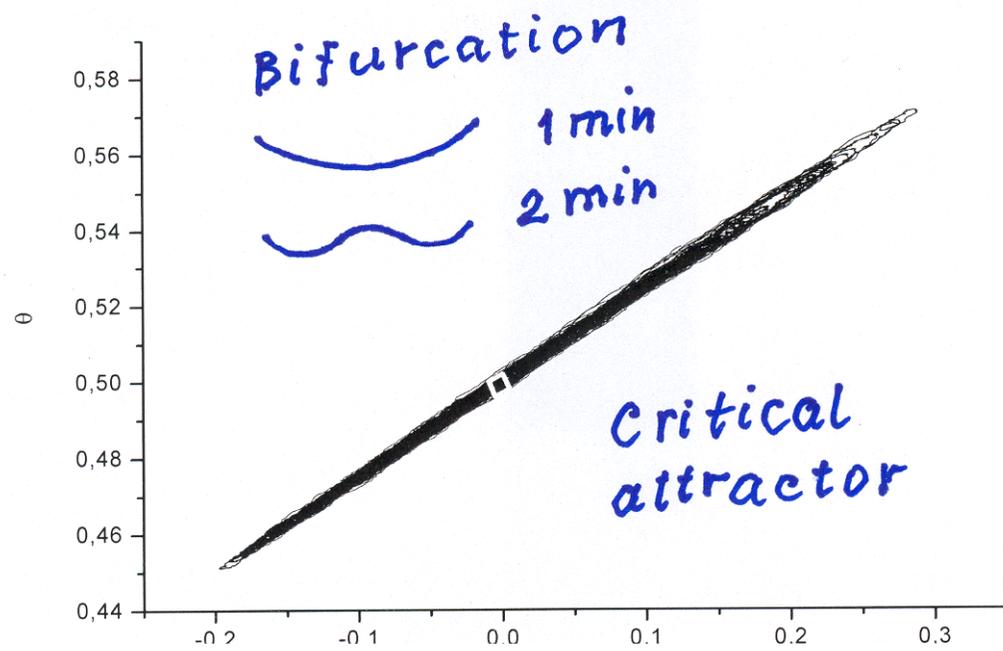
## Bifurcation:

distance between the limit point and Basin boundary decreases

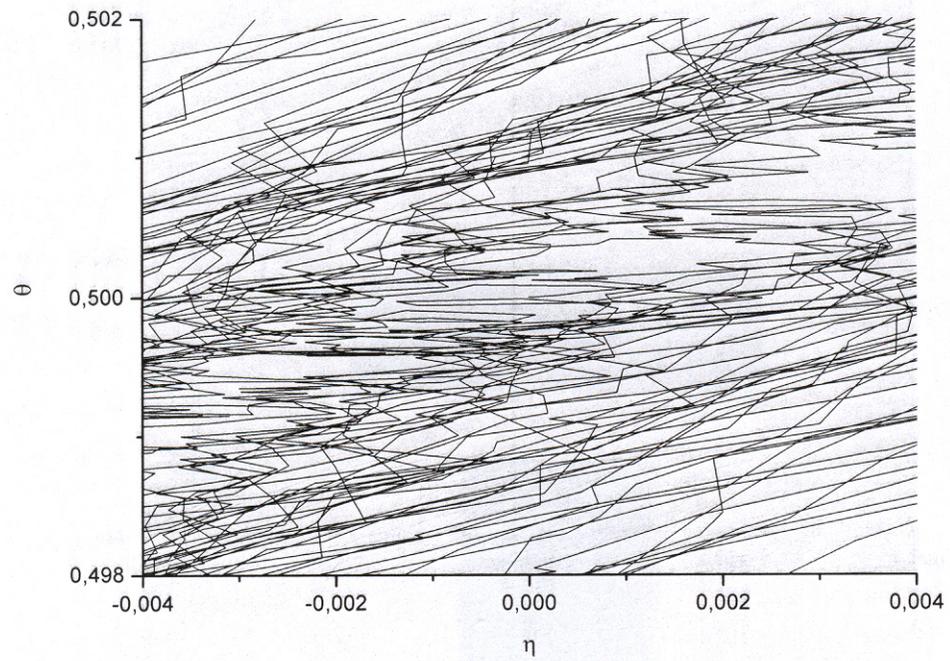
restoring force for deviations decreases

weak noise becomes strong

Langevin equations



# Brownian trajectory



В случае одной переменной уравнение для отклонений

$$dx/dt = -\lambda x + y(t).$$

$y(t)$

-- источник отклонений.

При начальном условии

$$x(t \rightarrow -\infty) = 0,$$

имеем

$$x(t) = e^{-\lambda t} \int_{-\infty}^t y(t') e^{\lambda t'} dt',$$

компонента Фурье

$$x_{\omega} = \frac{y_{\omega}}{\lambda - i\omega},$$

так что вещественная и мнимая части восприимчивости

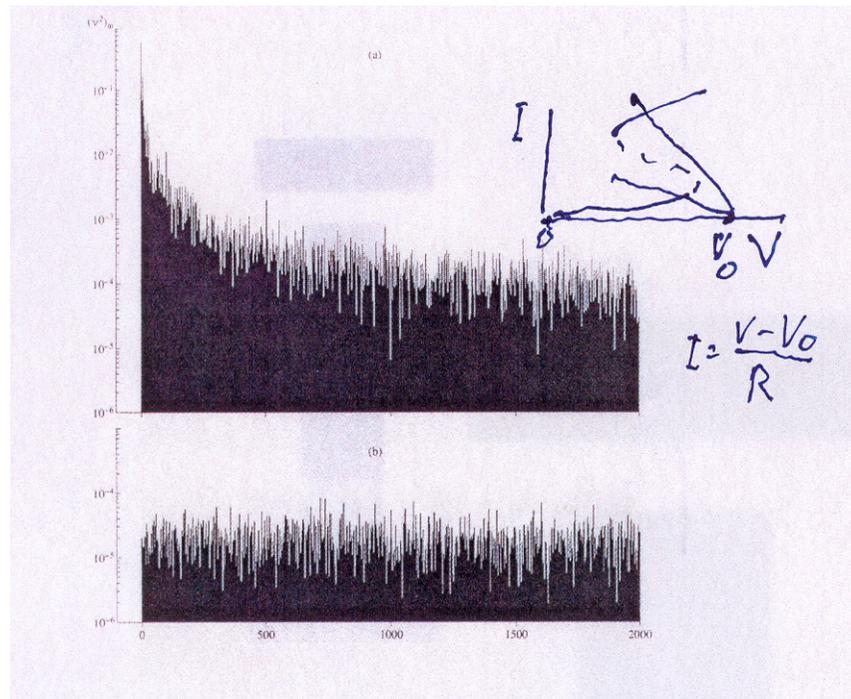
$$\operatorname{Re} \sigma = \frac{\lambda}{\lambda^2 - \omega^2}, \quad \operatorname{Im} \sigma = \frac{\omega}{\lambda^2 - \omega^2}.$$

Около бифуркации,

$$\lambda \rightarrow 0,$$

они расходятся на низких частотах,

$$\omega < \lambda.$$



$(v^2)_{\text{н.г}}$

